

# AN UNSTRUCTURED MESH NFT APPROACH TO ALL- SCALE ATMOSPHERIC FLOWS



**Mike Gillard, Joanna Szmelter**  
*Loughborough University*

## Mass, entropy and momentum conservation laws:

*Smolarkiewicz et. al. JCP 2016*

$$\frac{\partial \mathcal{G}_\varrho}{\partial t} + \nabla \cdot (\mathcal{G}_\varrho \mathbf{v}) = 0$$

$$\frac{\partial \mathcal{G}_\varrho \theta'}{\partial t} + \nabla \cdot (\mathcal{G}_\varrho \mathbf{v} \theta') = -\mathcal{G}_\varrho \left( \tilde{G}^T \mathbf{u} \cdot \nabla \theta_a - \mathcal{H} \right)$$

On a rotating sphere, cast in generalised time-dependent curvilinear coordinates, to enable a range of coordinate systems and 3D r-adaptivity.

$$\frac{\partial \mathcal{G}_\varrho \mathbf{u}}{\partial t} + \nabla \cdot (\mathcal{G}_\varrho \mathbf{v} \otimes \mathbf{u}) = -\mathcal{G}_\varrho \left( \Theta \tilde{G} \nabla \varphi + g \Upsilon_B \frac{\theta'}{\theta_b} + \mathbf{f} \times (\mathbf{u} - \Upsilon_C \mathbf{u}_a) - \mathcal{M}'(\mathbf{u}, \mathbf{u}, \Upsilon_C) - \mathcal{D} \right)$$

$$\varrho := \left[ \rho(\mathbf{x}, t), \frac{\rho_b(z) \theta_b(z)}{\theta(\mathbf{x}, t)}, \rho_b(z) \right] \quad \varphi := [c_p \theta_0 \pi', c_p \theta_0 \pi, c_p \theta_b \pi] \quad \Theta := \left[ \frac{\theta}{\theta_0}, \frac{\theta}{\theta_0}, 1 \right] \quad \Upsilon_B := \left[ \frac{\theta_b(z)}{\theta_a(\mathbf{x})}, \frac{\theta_b(z)}{\theta_a(\mathbf{x})}, 1 \right] \quad \Upsilon_C := \left[ \frac{\theta}{\theta_a(\mathbf{x})}, \frac{\theta}{\theta_a(\mathbf{x})}, 1 \right]$$

[compressible, pseudo-incompressible, anelastic]

## Advection: MPDATA

Multidimensional Positive Definite  
Advection Transport Algorithm

$$\frac{\partial G\Phi}{\partial t} + \nabla \cdot (\mathbf{V}\Phi) = G\mathcal{R}$$

$$\Phi_i^{n+1} = \mathcal{A}_i \left( \Phi^n + \frac{1}{2}\delta t\mathcal{R}, \mathbf{V}^{n+\frac{1}{2}}, G^n, G^{n+1} \right) + \frac{1}{2}\delta t\mathcal{R}^{n+1}$$

*Smolarkiewicz & Szmelter,  
JCP 2009, Acta Geo 2011*

## Helmholtz: from evolutionary form of the gas law

$$\begin{aligned} \frac{1}{\mathcal{G}} \nabla \cdot (\mathcal{G}\mathbf{v}^\nu) - \frac{1}{\xi} \frac{\phi_a}{\phi^{\nu-1}} \left( \frac{1}{\varrho^* \phi_a} \nabla \cdot (\varrho^* \phi_a \mathbf{v}^\nu) - \frac{1}{\varrho^*} \nabla \cdot (\varrho^* \mathbf{v}^\nu) \right) \\ - \frac{1}{\delta t \xi \phi^{\nu-1}} (\varphi^\nu - \hat{\varphi}) = 0 \end{aligned}$$

*Smolarkiewicz et. al. JCP 2016*

$$\Rightarrow \frac{1}{\zeta} \nabla \cdot \zeta \left( \check{\mathbf{v}} - \tilde{G}^T C \nabla \varphi \right)$$

## Implicit Integrator: GCR(k)

non-symmetric, preconditioned  
Krylov-subspace elliptic solver

For any initial guess,  $\phi^0$ , set  $r^0 = \mathcal{L}(\phi^0) - \mathcal{R}$ ,  $p^0 = \mathcal{P}^{-1}(r^0)$ ; then iterate:

For  $n = 1, 2, \dots$  until convergence

for  $\nu = 0, \dots, k-1$

$$\beta = \frac{\langle r^\nu \mathcal{L}(p^\nu) \rangle}{\langle \mathcal{L}(p^\nu) \mathcal{L}(p^\nu) \rangle},$$

$$\phi^{\nu+1} = \phi^\nu + \beta p^\nu,$$

$$r^{\nu+1} = r^\nu + \beta \mathcal{L}(p^\nu),$$

exit if  $\|r^{\nu+1}\| \leq \epsilon$ ,

$$e = \mathcal{P}^{-1}(r^{\nu+1}),$$

evaluate  $\mathcal{L}(e) = \frac{1}{\rho^*} \nabla \cdot C \nabla e$ ,

for  $l = 0, \dots, \nu$

$$\alpha_l = \frac{\langle \mathcal{L}(e) \mathcal{L}(p^l) \rangle}{\langle \mathcal{L}(p^l) \mathcal{L}(p^l) \rangle},$$

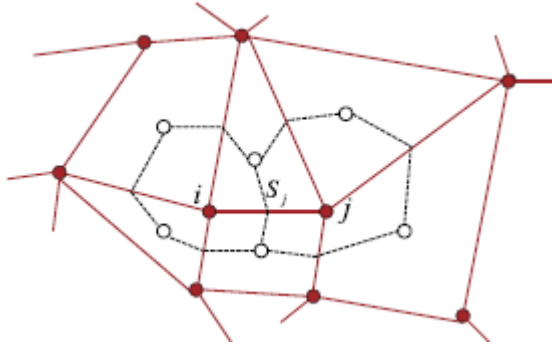
$$p^{\nu+1} = e + \sum_l^\nu \alpha_l p^l,$$

$$\mathcal{L}(p^{\nu+1}) = \mathcal{L}(e) + \sum_l^\nu \alpha_l \mathcal{L}(e^l),$$

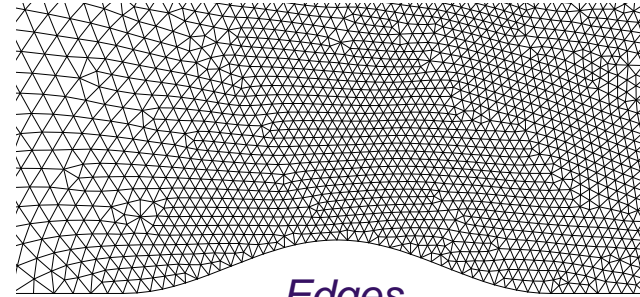
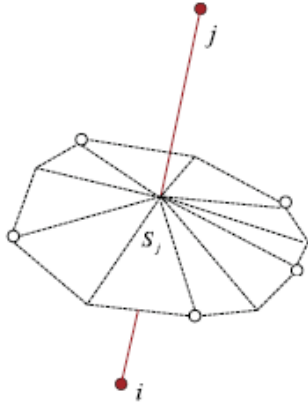
reset  $[\phi, r, e, \mathcal{L}(e)]^k$  to  $[\phi, r, e, \mathcal{L}(e)]^0$

*Smolarkiewicz &  
Margolin, 2000*

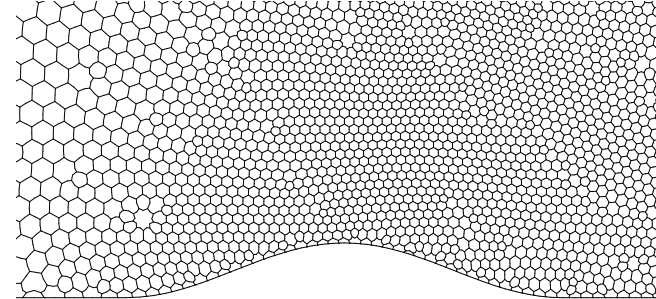
# Edge based finite volume discretisation



Uses a co-located data arrangement



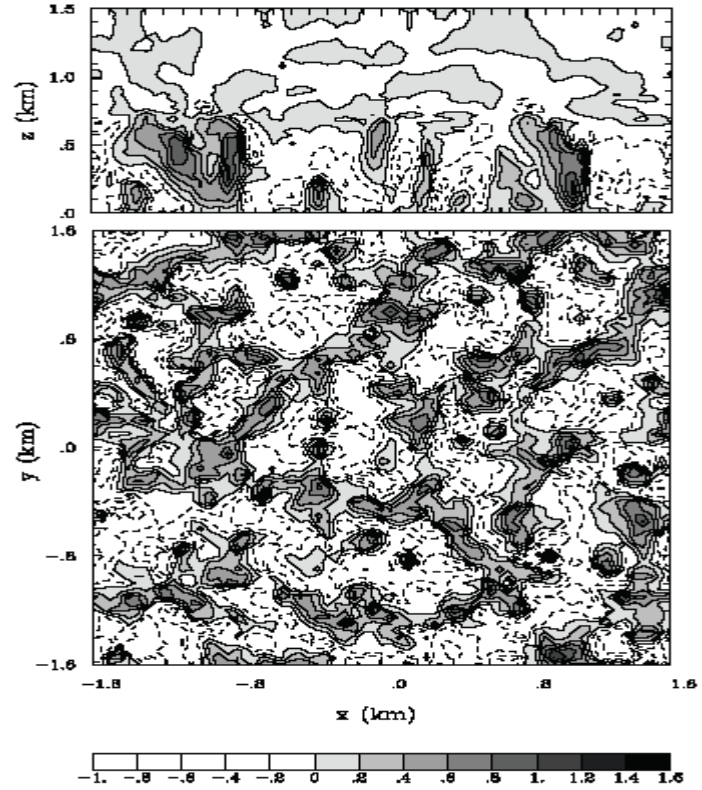
Edges



Median dual computational mesh  
Finite volumes

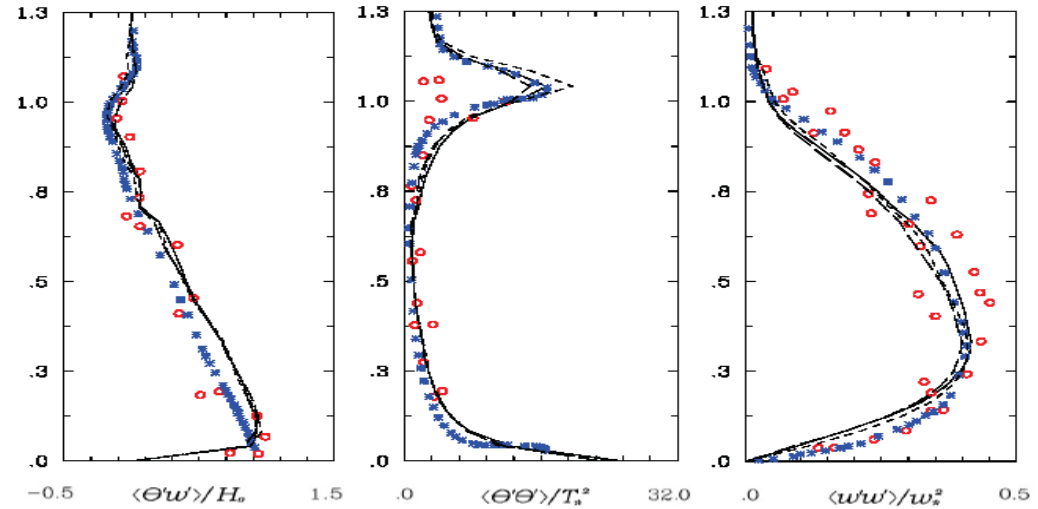
# Convective Planetary Boundary Layer

Smolarkiewicz et al JCP 2013

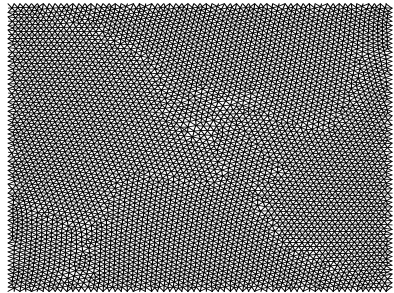


Vertical velocity in central xz cross section (top) and the horizontal plane (bottom): instantaneous solution for triangular prismatic mesh after ~13 eddy turnover times

Schmidt & Schumann JFM 1989



Vertical profile of dimensionless resolved heat flux, temperature variance, and vertical velocity; for dimensionless height  $z/z_i$



- Edge-based T ———
- Edge-based C - - -
- EULAG - · - · -
- SS LES \*
- Observation ○

# Preconditioning

Elementary preconditioner,  
slower elliptic solver

Advanced Preconditioner  
faster elliptic solver

$$I \quad \leftarrow \quad \mathcal{P} \quad \rightarrow \quad \mathcal{L}$$

Choose  $\mathcal{P}$  close to  $\mathcal{L}$ , but easier to solve.

Take only diagonal terms in  $\mathcal{C}$ :  $\mathcal{P} \approx \mathcal{L}$

Invert  $\mathcal{P}_z$  implicitly, with  $\mathcal{P}_H$  lagged:

$$\frac{\partial e}{\partial \tilde{\tau}} = \mathcal{P}(e) - r \quad \Rightarrow \quad \frac{ee^{\mu+1} - ee^{\mu}}{\Delta \tilde{\tau}} = \mathcal{P}_H(e^{\mu}) + \mathcal{P}_z(e^{\mu+1}) - r$$

$$(I - \Delta \tilde{\tau} \mathcal{P}_z) e^{\mu+1} = e^{\mu} + \Delta \tilde{\tau} (\mathcal{P}_H(e^{\mu}) - r^{\nu+1})$$

For any initial guess,  $\phi^0$ , set  $r^0 = \mathcal{L}(\phi^0) - R$ ,  $p^0 = \mathcal{P}^{-1}(r^0)$ ; then iterate:

For  $n = 1, 2, \dots$  until convergence

for  $\nu = 0, \dots, k - 1$

$$\beta = \frac{\langle r^{\nu} \mathcal{L}(p^{\nu}) \rangle}{\langle \mathcal{L}(p^{\nu}) \mathcal{L}(p^{\nu}) \rangle},$$

$$\phi^{\nu+1} = \phi^{\nu} + \beta p^{\nu},$$

$$r^{\nu+1} = r^{\nu} + \beta \mathcal{L}(p^{\nu}),$$

exit if  $\|r^{\nu+1}\| \leq \epsilon$ ,

$$e = \mathcal{P}^{-1}(r^{\nu+1}),$$

$$\text{evaluate } \mathcal{L}(e) = \frac{1}{\rho^*} \nabla \cdot C \nabla e,$$

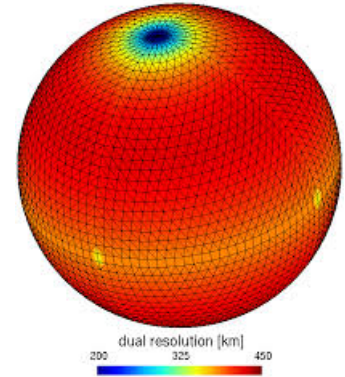
for  $l = 0, \dots, \nu$

$$\alpha_l = \frac{\langle \mathcal{L}(e) \mathcal{L}(p^l) \rangle}{\langle \mathcal{L}(p^l) \mathcal{L}(p^l) \rangle},$$

$$p^{\nu+1} = e + \sum_l^{\nu} \alpha_l p^l,$$

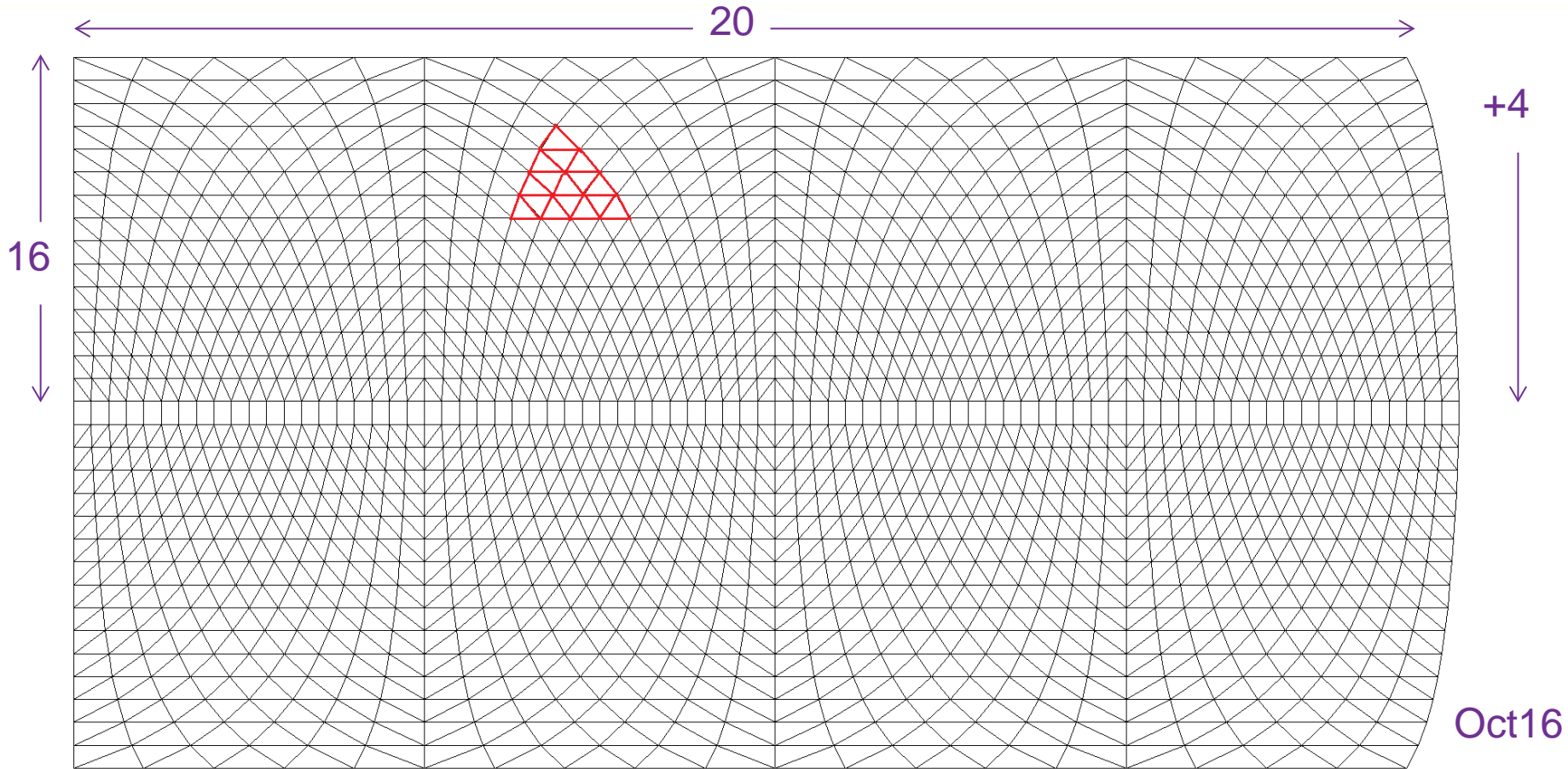
$$\mathcal{L}(p^{\nu+1}) = \mathcal{L}(e) + \sum_l^{\nu} \alpha_l \mathcal{L}(e^l),$$

reset  $[\phi, r, e, \mathcal{L}(e)]^k$  to  $[\phi, r, e, \mathcal{L}(e)]^0$





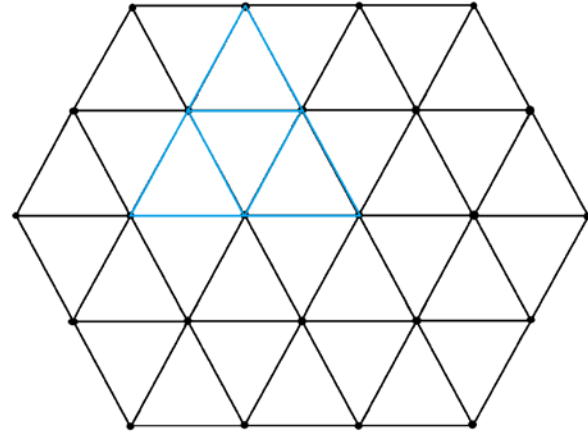
# Multigrid meshes using Atlas



# Multigrid meshes using Atlas

Octahedral mesh has  
reasonably consistent  
structure

Latitude  $n = 1 \dots N$   
has  $20 + 4n$  longitudes





# Multigrid meshes using Atlas

Remove every other latitude

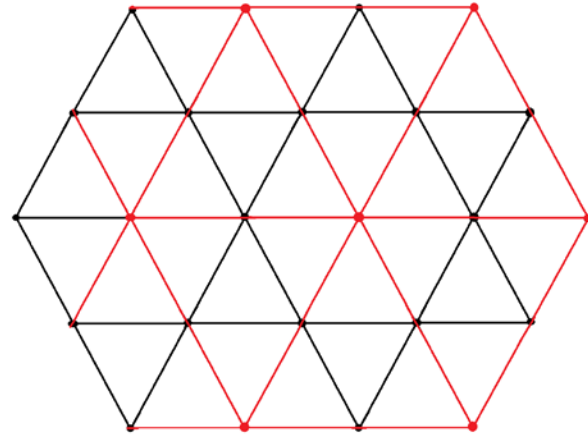
Half the number of longitudes at the remaining latitudes

Reduces horizontal nodes by a factor of  $\sim 4$

Need to pick starting latitude

Coarse mesh is almost octahedral

All coarse mesh nodes coincide with nodes on the fine mesh



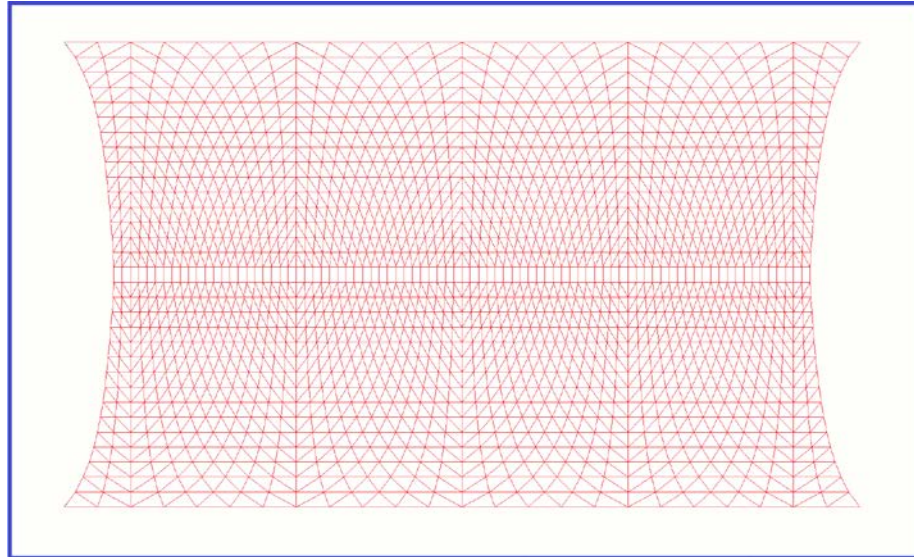
# Multigrid meshes using Atlas

Octahedral 16 mesh:

Remove odd latitudes

computational domain

Single level of mesh coarsening



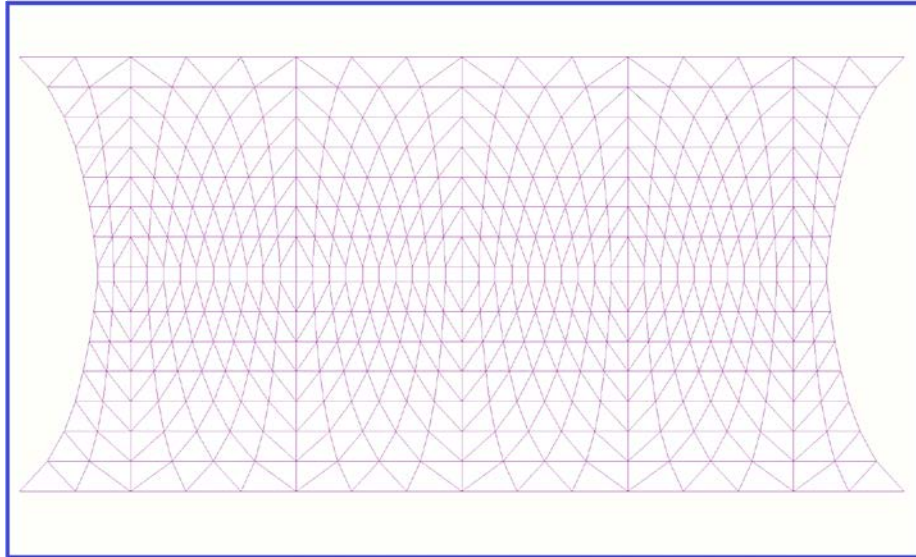
# Multigrid meshes using Atlas

Octahedral 16 mesh:

Remove odd latitudes

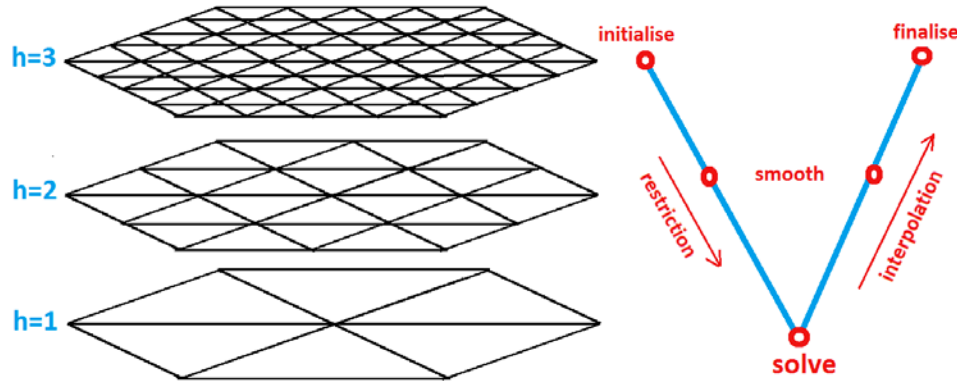
computational domain

Single level of mesh coarsening



# Multigrid Preconditioning

## V-Cycle



Smoother / Solver, varying number of iterations depending requirements

$$(I - \Delta \tilde{\tau} \mathcal{P}_z) e^{\mu+1} = e^{\mu} + \Delta \tilde{\tau} (\mathcal{P}_H(e^{\mu}) - r^{\nu+1})$$

V-Cycle: Solve on  $h=1$ , else  $\text{cycle}(e_h, r_h)$

$\text{cycle}(e_h, r_h)$ :

Smooth( $e_h$ )

calculate residule:  $rr_h = \mathcal{P}(e_h) - r_h$

restrict  $rr_h \rightarrow rr_{h-1}$ :  $rr_{h-1} = \mathcal{I}_h^{h-1} rr_h$

initialise solution error:  $ee_{h-1} [\equiv e - e']$

$\text{cycle}(ee_{h-1}, rr_{h-1})$

Interpolate solution error:  $ee_h = \mathcal{I}_{h-1}^h ee_{h-1}$

Add coarse grid error:  $e_h = e_h - ee_h$

Smooth( $e_h$ )

# Multigrid Preconditioning

Baroclinic Instability – 9 days.

Compressible Nonhydrostatic Equations.

Multigrid / implicit vertical preconditioner.

Multigrid:

Wall time ~ 4hrs

Elliptic solver iterations average – 3:0

Without multigrid:

Wall time ~ 23hrs

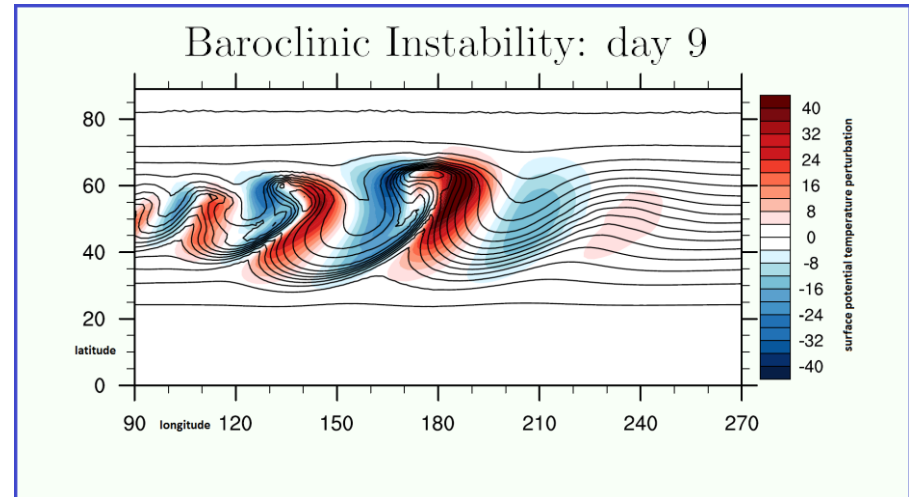
Elliptic solver iterations average – 60:7

Setup:

Octahedral 80 grid,

time-step 900s,

1/1 task/thread,



- NFT-MPDATA solvers based on an edge based finite volume method for unstructured grids provide a high quality results for a wide range of mesh types and is applicable for all scales atmospheric flows.
- Substantial efficiency gains have been achieved by introducing a horizontal multigrid preconditioner to a Krylov solver of Helmholtz equations forming a part of a global nonhydrostatic model.

